Effective dynamics for low-amplitude transient elastic waves in a 1D periodic array of non-linear interfaces

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Abstract

This presentation focuses on the time-domain propagation of elastic waves through a 1D periodic medium that contains non-linear imperfect interfaces, i.e. interfaces exhibiting a discontinuity in displacement and stress governed by a non-linear constitutive relation. In this context, we investigate transient waves with both lowamplitude and long-wavelength, and aim at deriving homogenized models that describe their effective motion.

Keywords: Homogenization – Correctors – Imperfect interfaces – Non-linear waves – Timedomain numerical simulations

1 Objectives

The array of interfaces considered is generated by a, possibly heterogeneous, cell repeated periodically and bonded by interfaces that are associated with transmission conditions of nonlinear "spring-mass" type. More precisely, the imperfect interfaces are characterized by a linear dynamics but a non-linear elasticity law. The latter is not specified at first and only key theoretical assumptions are required. To establish an effective model, the two-scale asymptotic homogenization method is deployed, up to the first-order. To begin, an effective model is obtained for the leading zeroth-order contribution to the microstructured wavefield. It amounts to a wave equation with a non-linear constitutive stress-strain relation that is inherited from the behavior of the imperfect interfaces at the microscale. The next first-order corrector term is then shown to be expressed in terms of a cell function and the solution of a linear elastic wave equation. Without further hypothesis, the constitutive relation and the source term of the latter depend non-linearly on the zeroth-order field, as does the cell function. Combining these zeroth- and first-order models leads to an approximation of both the macroscopic behavior of the microstructured wavefield and its smallscale fluctuations within the periodic array.

2 Setting: microstructured configuration

We consider the propagation of transient waves in a 1D periodic elastic medium containing imperfect interfaces. The latter have spacing hand, for simplicity but with no loss of generality, we consider that they are located at $X_n = nh$ with $n \in \mathbb{Z}$. The elastic medium is supposed to be h-periodic and linear elastic with mass density $\rho_h(X)$ and Young's modulus $E_h(X)$. Given a source term F, the displacement field U_h is governed by the time-domain wave equation

$$\rho_h(X)\frac{\partial^2 U_h}{\partial t^2}(X,t) = \frac{\partial \Sigma_h}{\partial X}(X,t) + F(X,t) \quad (1)$$

where

$$\Sigma_h(X,t) = E_h(X) \frac{\partial U_h}{\partial X}(X,t)$$

with Σ_h being the stress field. Moreover, the interfaces are assumed to be characterized by the interface mass and rigidity parameters Mand K, respectively, together with the, possibly non-linear, constitutive relation \mathcal{R} , so that the following transmission conditions apply at any interface point X_n , see [1–4]:

$$\begin{cases} M \left\langle\!\!\left\langle \frac{\partial^2 U_h}{\partial t^2}(\cdot, t) \right\rangle\!\!\right\rangle_{X_n} = \llbracket \Sigma_h(\cdot, t) \rrbracket_{X_n} & \text{(2a)} \\ \left\langle\!\!\left\langle \Sigma_h(\cdot, t) \right\rangle\!\!\right\rangle_{X_n} = K \mathcal{R} \left(\llbracket U_h(\cdot, t) \rrbracket_{X_n}\right), \text{(2b)} \end{cases}$$

where, for any function g(X), we define the jump and mean operators $\llbracket \cdot \rrbracket_{X_n}$ and $\langle \! \langle \cdot \rangle \! \rangle_{X_n}$ as

$$[[g]]_{X_n} = g(X_n^+) - g(X_n^-),$$

$$\langle\!\langle g \rangle\!\rangle_{X_n} = \frac{1}{2} \big(g(X_n^+) + g(X_n^-) \big).$$
 (3)

In addition, both the displacement U_h and the stress field Σ_h are continuous on the open intervals (X_n, X_{n+1}) .

3 Main homogenization results

We now consider a reference wavelength λ^* and introduce the following parameters

$$k^* = \frac{2\pi}{\lambda^*}$$
 and $\eta = hk^*$, (4)

 k^* being the reference wavenumber. In this study it is assumed that $\eta \ll 1$ and that the source term F is of relatively *low-amplitude* (an issue that will be discussed). The objective is to derive an effective dynamical model, up to the first-order, for the waves propagating in the periodic interface array considered. More precisely, we seek an approximation $U^{(1)}$ of the solution U_h to (1–2) of the form:

$$U_h(X,t) = U^{(1)}(X,t) + o(h).$$

The main results of this study is that the soughtafter approximation is given by

$$U^{(1)}(X,t) = U_0(X,t) + hU_1(X,t), \quad (5)$$

where the zeroth-order field U_0 in (5) is continuous and is solution of the problem

$$\rho_{\text{eff}} \frac{\partial^2 U_0}{\partial t^2}(X, t) = \frac{\partial \Sigma_0}{\partial X}(X, t) + F(X, t)$$

with

$$\Sigma_0(X,t) = \mathcal{G}_{\text{eff}}\big(\mathcal{E}_0(X,t)\big).$$

Here, $\mathcal{E}_0 = \partial U_0 / \partial X$ and \mathcal{G}_{eff} is an effective strainstress relation that is local and, generally speaking, non-linear, while ρ_{eff} is an effective mass density. Moreover, the first-order corrector field U_1 in (5) can be written as

$$U_1(X,t) = \overline{U}_1(X,t) + \mathcal{P}(y,\mathcal{E}_0(X,t)) \mathcal{E}_0(X,t)$$

with y = (X - nh)/h for $X \in (nh, (n+1)h)$ and where the *cell function* \mathcal{P} is, generally speaking, a non-linear function of \mathcal{E}_0 . The *mean* field \overline{U}_1 is solution to the *linear* problem:

$$\rho_{\text{eff}} \frac{\partial^2 \overline{U}_1}{\partial t^2}(X, t) = \frac{\partial \overline{\Sigma}_1}{\partial X}(X, t) + \mathcal{S}\big(U_0(X, t)\big)$$

with

$$\overline{\Sigma}_1(X,t) = \mathcal{G}'_{\text{eff}}(\mathcal{E}_0(X,t)) \overline{\mathcal{E}}_1(X,t),$$

where $\overline{\mathcal{E}}_1 = \partial \overline{U}_1 / \partial X$, while both the parameter $\mathcal{G}'_{\text{eff}}(\mathcal{E}_0(X,t))$, which is the derivative of \mathcal{G}_{eff} , and the source term $\mathcal{S}(U_0(X,t))$ depend explicitly on the zeroth-order field, locally in space and time, and in a non-linear fashion.

Particularizing for a prototypical non-linear interface law and in the cases of a homogeneous periodic cell and a bilaminated one, the behavior of the obtained models will then be illustrated on a set of numerical examples and compared with full-field simulations. Both the influence of the dominant wavelength and of the wavefield amplitude will be investigated numerically, as well as the characteristic features related to non-linear phenomena.

References

- J. D. Achenbach and A. N. Norris. Loss of specular reflection due to nonlinear crackface interaction. *Journal of Nondestructive Evaluation*, 3(4):229–239, 1982.
- [2] S.C. Bandis, A.C. Lumsden, and N.R. Barton. Fundamentals of rock joint deformation. International Journal of Rock Mechanics and Mining Sciences & Geomechanics Abstracts, 20(6):249 – 268, 1983.
- [3] D. Broda, W.J. Staszewski, A. Martowicz, T. Uhl, and V.V. Silberschmidt. Modelling of nonlinear crack-wave interactions for damage detection based on ultrasound—a review. *Journal of Sound and Vibration*, 333(4):1097 – 1118, 2014.
- [4] I. Sevostianov, R. Rodriguez-Ramos, R. Guinovart-Diaz, J. Bravo-Castillero, and F.J. Sabina. Connections between different models describing imperfect interfaces in periodic fiber-reinforced composites. *International Journal of Solids* and Structures, 49(13):1518 – 1525, 2012.